Blackbody Radiation

Let us use Einstein's approach to relate gain and spontaneous emission.

According to Plank:

1.The probability of radiation taking place from a black body decreased as the frequency of the radiation increased.

2. The black body radiation flow out in discrete, lumpy quantities *hv*

MODEL



Planck postulated:

$$\left< S \right>_{hv} \propto \frac{(hv)^2}{exp\left(\frac{hv}{kT}\right) - 1}$$

$$h = 6.626 \cdot 10^{-34} J \cdot s$$

Optical gain spectrum measurements: From TSE

Einstein's approach to gain and spontaneous emission

Change of the photon numbers S:

$$\Delta S = -B_{STIM}^{ABS} \cdot S \cdot f_1 \cdot (1 - f_2) + B_{STIM}^{EM} \cdot S \cdot f_2 \cdot (1 - f_1) + A_{SPON}^{EM} \cdot f_2 \cdot (1 - f_1)$$

Possible contribution of the various electron transitions to emission/absorption of photons in energy interval [hv,hv+dhv] is accounted for in A and B constants.



Comments



Optical gain spectrum measurements: <u>From TSE</u> <u>Gain and spontaneous emission relation</u>

No equilibrium: $E_{F1} \neq E_{F2}, \Delta S \neq 0$.



$$G(h\nu) = \frac{\Delta S_{STIM}}{\Delta x} \cdot \frac{1}{S} = \frac{B}{\Delta x} \cdot (f_2 - f_1)$$

$$R_{SPON}(h\nu) = \frac{\Delta S_{SPON}}{\Delta t} = \upsilon_g \cdot \frac{\Delta S_{SPON}}{\Delta x} = \upsilon_g \cdot \frac{A}{\Delta x} \cdot f_2 \cdot (1 - f_1)$$

$$R_{SPON}(h\nu) \propto G(h\nu) \cdot (h\nu)^2 \cdot \frac{f_2 \cdot (1 - f_1)}{(f_2 - f_1)}$$

$$I_{SP}(hv) \propto \frac{G(hv) \cdot (hv)^2}{\left(1 - \exp\left(\frac{hv - \Delta E_F}{kT}\right)\right)}$$

Comments



Optical gain spectrum measurements: From TSE





Determination of transparency energy: Andrekson technique



NO VOLTAGE CHANGE ACROSS DIODE WHEN $h\nu = \Delta E_F$!

Determination of transparency energy: Andrekson technique



The transparency energy is equal to $\Delta E_F = (E_e)_F - (E_h)_F$.

* Induced voltage is "in phase" with ASE for $hv > (E_e)_F - (E_h)_F$ since larger absorption means larger voltage. For $hv < (E_e)_F - (E_h)_F$ modes supported by resonator are amplified stronger and maximum of ASE corresponds to the voltage minimum. The energy of voltage oscillation "phase change" corresponds to transparency energy.

Measurements of optical loss by Variable cavity length method



 $\frac{1}{\eta_{\text{ext}}} = \frac{1}{\eta_{\text{i}}} \cdot \left(1 + \frac{\alpha_{\text{i}}}{\alpha_{\text{m}}}\right)$ $\alpha_{\text{m}} = \frac{1}{L} \cdot \ln \frac{1}{\sqrt{R_{1}R_{2}}}$ Assumption: η_{i} and α_{i} are cavity length

 $\eta_{\text{ext}} = \eta_{\text{i}} \cdot \frac{\alpha_{\text{m}}}{\alpha_{\text{m}} + \alpha_{\text{i}}}$

independent

Advantage: Simplicity of measurements

Shortcoming: For short cavities or small internal loss threshold is L-dependent and errors arise. Also a lot of laser material is required.

Measurements of optical loss from modal gain spectrum



$$g^{\text{TE/TM}}(h\nu) = \Gamma^{\text{TE/TM}} \cdot G^{\text{TE/TM}}(h\nu) - \alpha_{\text{tot}}^{\text{TE/TM}}$$



1. TE and TM modal gains intersection

TE – comes from C-HH transition; TM – comes from C-LH transition. Thus, spectra for TE and TM gains are different and corresponding gains and can be equal only when **material** gain G=0 (transparency point).

<u>2. Saturation for hv < E_g</u>

For photon energies below bandgap, material gain is equal to zero. Modal (g) gain is equal to total loss within this spectra region. Usually, long wavelength tail of the modal gain spectra gives total loss value. When loss are determined, transparency energy (condition G=0) can be estimated from the gain spectrum.